

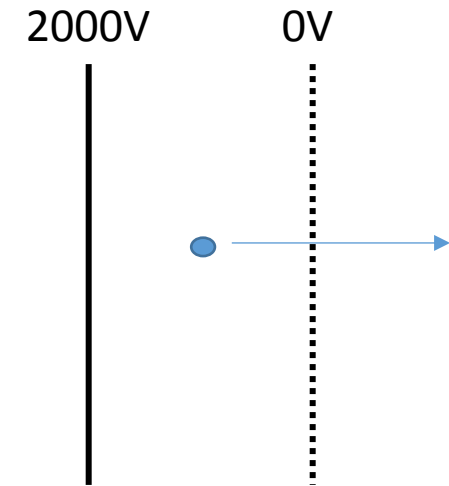
Time-of-flight mass spectrometer (TOF)

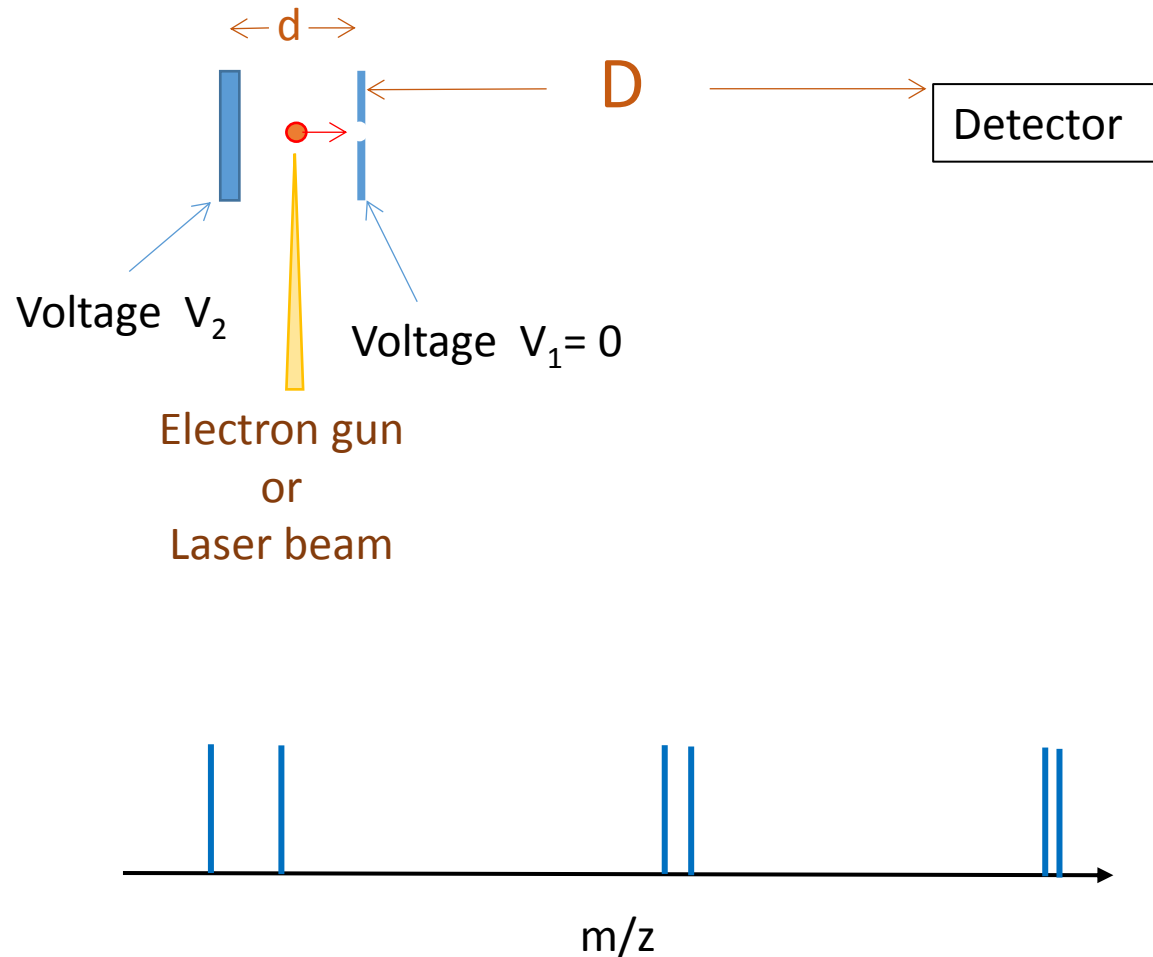
Velocity gained from electric field

$$E = 0.5 \text{ mV}^2$$

$$1000 \text{ eV} \times 1.6 \times 10^{-12} \text{ erg/eV} = 0.5 \times [100 / (6 \times 10^{23})] \times V^2$$

$$V = 4.4 \times 10^6 \text{ cm/s}$$





Time to arrive detector

$$T = t_1 + t_2$$

$$0.5xd = V_0 t_1 + 0.5x(F/m)x(t_1)^2$$

$$d = [qx\hat{E}/m]x(t_1)^2 \quad (\text{assume } V_0 = 0)$$

$$t_1 = \text{sqrt}[dxm/(qx\hat{E})]$$

$$t_2 = D/V$$

$$V = V_0 + axt_1$$

$$= qx\hat{E}/mxt_1$$

$$= \text{sqrt}[(qx\hat{E})/m]$$

$$t_2 = Dx\text{sqrt}[m/(qx\hat{E})]$$

$$\text{Arrival time: } T = t_1 + t_2 \sim \text{sqrt}(m)$$

Example 1

$$d = 2 \text{ cm} = 0.02 \text{ m}$$

$$D = 1 \text{ m}$$

$$V1 = 2000 \text{ V}$$

For $m=100 \text{ Da}$

$$\begin{aligned} t1 &= \text{sqrt}[0.02 \times 100 \times 10^{-3} / (6.02 \times 10^{23} \times 1.6 \times 10^{-19} \times 2000)] \\ &= \text{sqrt}(10^{-9}) = 3.162 \times 10^{-5} \text{ s} \end{aligned}$$

$$\begin{aligned} t2 &= 1 \times \text{sqrt}[100 \times 10^{-3} / (6.02 \times 10^{23} \times 1.6 \times 10^{-19} \times 2000)] \\ &= \text{sqrt}(5 \times 10^{-8}) = 2.236 \times 10^{-4} \text{ s} \end{aligned}$$

$$T = 255.22 \text{ } \mu\text{s}$$

For $m=101 \text{ Da}$

$$t1 = 3.178 \times 10^{-5} \text{ s}$$

$$t2 = 2.247 \times 10^{-4} \text{ s}$$

$$T = 256.48 \text{ } \mu\text{s}$$

$$\Delta T = 1.26 \text{ } \mu\text{s}$$

For $m= 900 \text{ Da}$

$$t1 = 9.486 \times 10^{-5} \text{ s}$$

$$t2 = 6.708 \times 10^{-4} \text{ s}$$

$$T = 765.66 \text{ } \mu\text{s}$$

Scan from $m=0-900 \text{ Da}$ takes $765 \text{ } \mu\text{s}$

For $m= 901 \text{ Da}$

$$t1 = 9.492 \times 10^{-5} \text{ s}$$

$$t2 = 6.712 \times 10^{-4} \text{ s}$$

$$T = 7656.12 \text{ } \mu\text{s}$$

$$\Delta T = 0.5 \text{ } \mu\text{s}$$

Example 2

$$d = 2 \text{ cm} = 0.02 \text{ m}$$

$$D = 2 \text{ m}$$

$$V1 = 2000 \text{ V}$$

For $m=100 \text{ Da}$ and $m=101 \text{ Da}$

$$\Delta T = 1.26 \times \sqrt{2} \text{ } \mu\text{s} = 1.78 \text{ } \mu\text{s}$$

Time-of-flight mass spectrometer (TOF)

- No mass range (High mass range is limited by detector (detector has small response to slow particles))
- High sensitivity: most ions generated by ionization arrive detector
- Fast scan
- Easy construction

Distance deviated by magnetic field

$$S = V_0 t + 0.5 a t^2$$

$$= 0.5 F / m (D / v)^2$$

$$= 0.5 q v B / m (D / v)^2$$

$$= 0.5 q B D^2 / (m v)$$

$$= 0.5 q B D^2 / \sqrt{2 E m} \quad \text{where } E \text{ is the kinetic energy}$$

For large mass, the deviation is small.

Magnetic field has larger effect on electrons than ions

Effect of earth magnetic field

For an single charged ion, $q = 1.6 \times 10^{-19} \text{ C}$

Assume $m = 100 \times 10^{-3} / (6 \times 10^{23}) \text{ kg}$, kinetic energy = 100 eV

$$100 \text{ eV} \times 1.6 \times 10^{-19} \text{ J/eV} = 0.5 \times 100 \times 10^{-3} \text{ kg} / (6.02 \times 10^{23}) \times v^2$$
$$v = 13000 \text{ m/s}$$

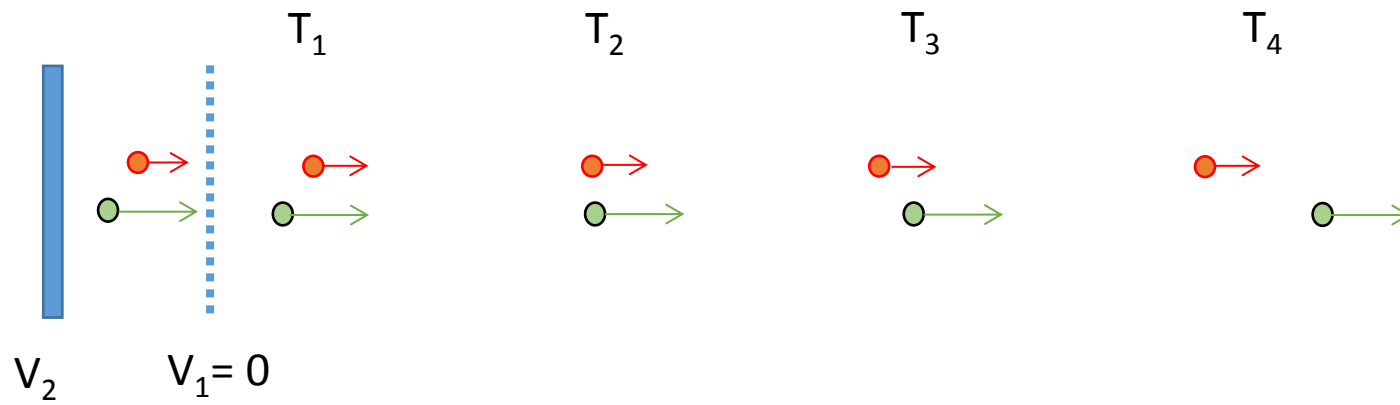
Earth magnetic field: $20\text{-}60 \times 10^{-6} \text{ tesla}$

$$F = 1.6 \times 10^{-19} \text{ C} \times 13000 \text{ m/s} \times 50 \times 10^{-6} \text{ tesla} = 1.04 \times 10^{-19} \text{ N}$$

Flight distance 1 m, time $t = 1/13000 = 7.7 \times 10^{-5} \text{ s}$

$$S = V_0 \times t + 0.5 \times a \times t^2$$
$$= 0.5 \times 1.04 \times 10^{-19} \text{ N} / [100 \times 10^{-3} / (6 \times 10^{23})] \times (7.7 \times 10^{-5})^2$$
$$= 1.8 \times 10^{-3} \text{ m}$$

Effects of different initial positions



Time-of-Flight Mass Spectrometer with Improved Resolution

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(Received June 30, 1955; revised version received August 22, 1955)

A new type of ion gun is described which greatly improves the resolution of a nonmagnetic time-of-flight mass spectrometer. The focusing action of this gun is discussed and analyzed mathematically. The validity of the analysis and the practicability of the gun are demonstrated by the spectra obtained. The spectrometer is capable of measuring the relative abundance of adjacent masses well beyond 100 amu.

INTRODUCTION

THIS paper describes the improved mass resolution made possible by a new ion gun in a nonmagnetic, time-of-flight (TOF) mass spectrometer. Although the properties of this gun may make it useful for focusing electrically accelerated ion beams in other applications, the discussion will be restricted to mass spectrometers which have been tested.

Following a brief description of TOF spectrometers, their resolution is discussed. The new ion gun is then described and its focusing action is analyzed mathematically. Experimental results substantiating the general theoretical conclusions are presented. Finally, a

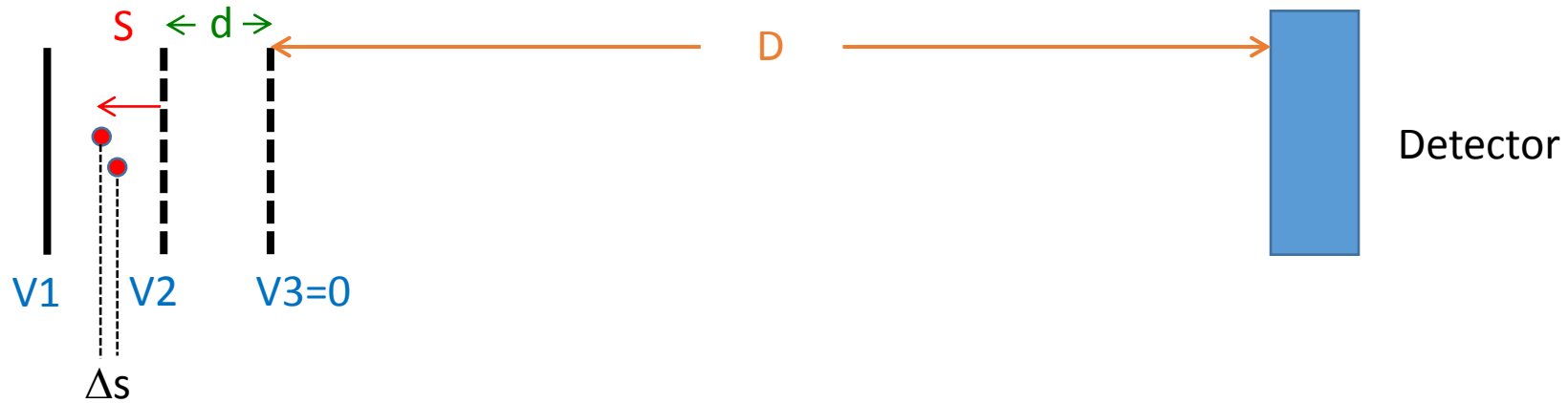
they have separated into bunches corresponding to q/m . If only singly charged ions are present, the lightest group reaches the detector first and is followed by groups of successively heavier mass. Thus, each source pulse results in a mass spectrum which can easily be displayed by connecting the ion collector to the vertical plates of an oscilloscope. An expanded portion of such a spectrum (obtained with the improved ion gun) is shown in Fig. 1.

The known abundances of the xenon isotopes displayed are 128, 2%; 129, 26.4%; 130, 4.1%; 131, 21.2%; 132, 26.9%; 134, 10.4%; 136, 8.9%.

Wiley-Maclaren ToFMS

$$T = t_1(V_1, V_2, s) + t_2(V_1, V_2, d) + t_3(V_1, V_2, D)$$

For a given s, d, D , find the V_1 and V_2 such that $\frac{\partial T}{\partial \Delta s} = 0$



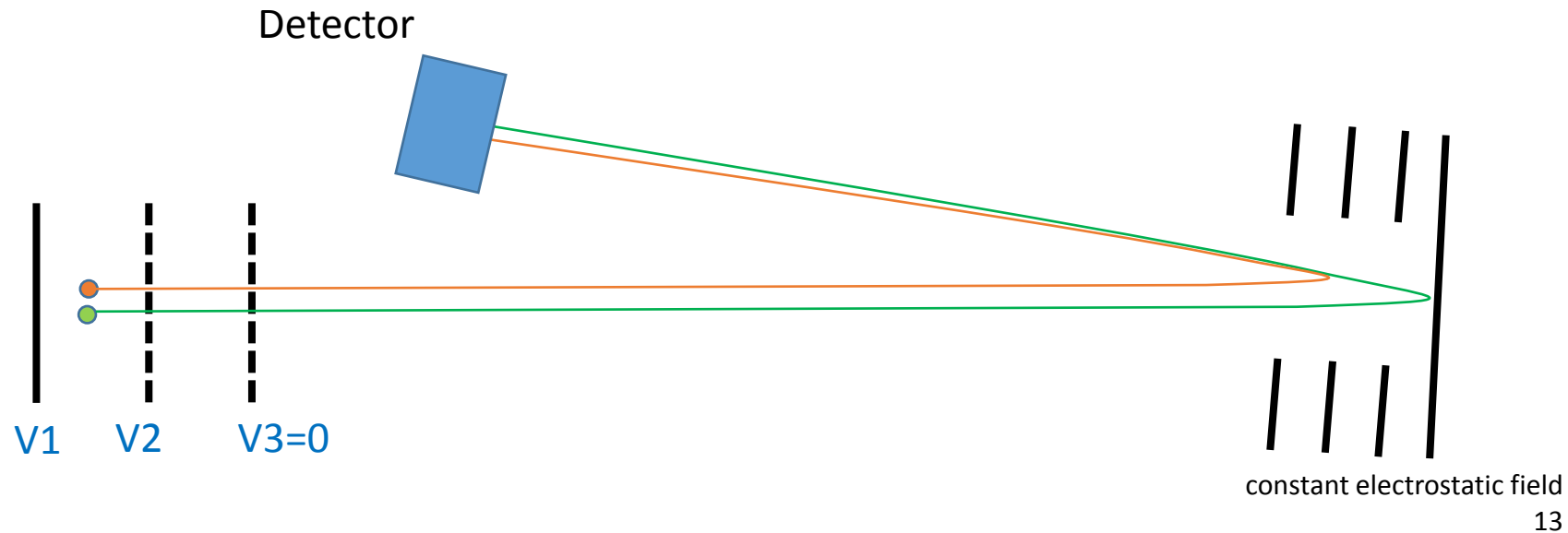
Delayed extraction

Ions generated from MALDI have different velocities. Ions with low velocity locate at position further away from detector, but they get more energy from electric field. The energy turns into velocity and they catch up the other ions at the detector.



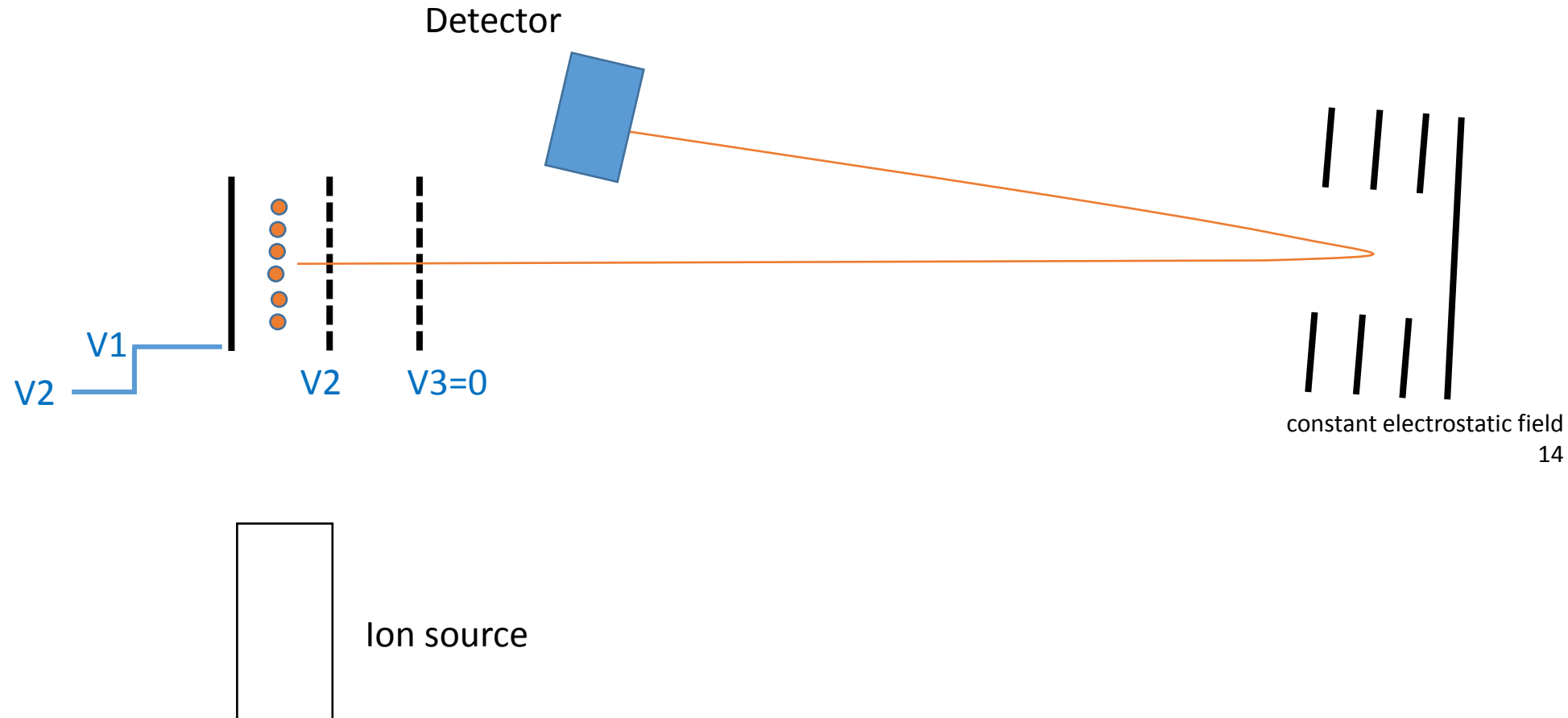
Reflectron TOF

The kinetic energy distribution in the direction of ion flight can be corrected by using a reflectron. The more energetic ions penetrate deeper into the reflectron, and take a slightly longer path to the detector. Less energetic ions of the same mass-to-charge ratio penetrate a shorter distance into the reflectron and, correspondingly, take a shorter path to the detector. The flat surface of the ion detector (typically a microchannel plate, MCP) is placed at the plane where ions of same m/z but with different energies arrive. They arrive the detector at the same time.



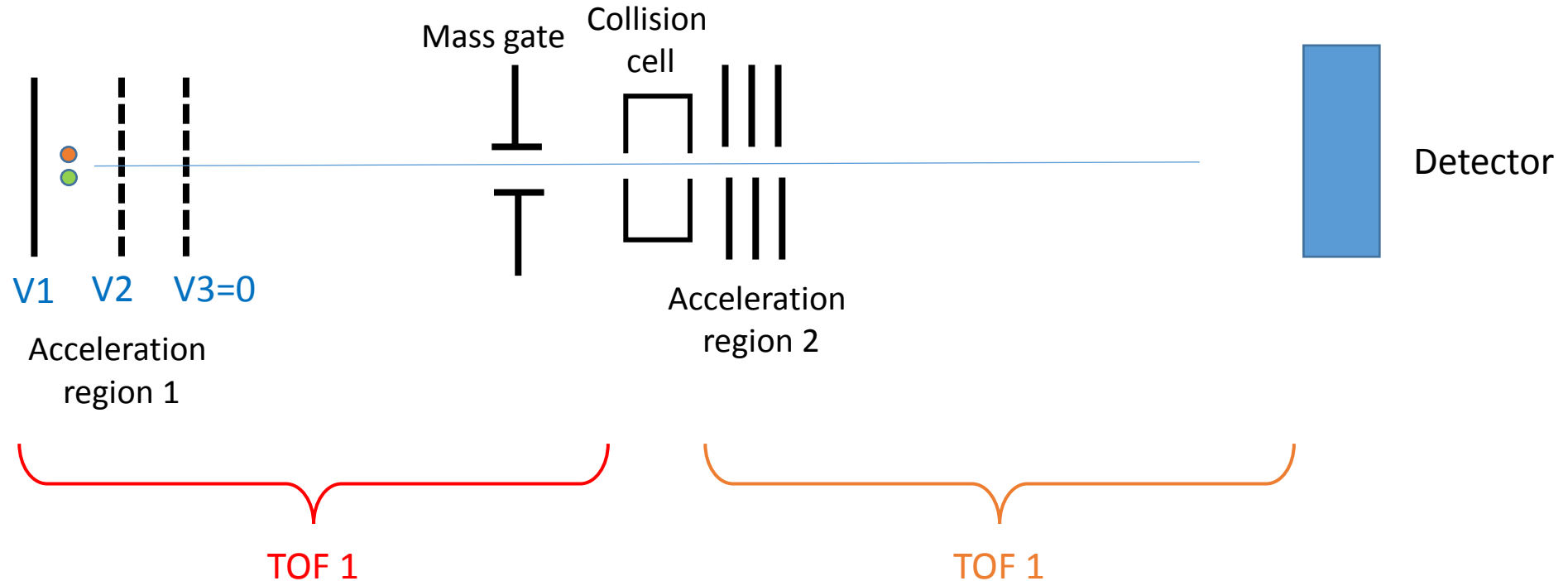
Orthogonal acceleration time-of-flight

Ions from continuous ion sources are injected into a TOF by "orthogonal extraction" in which ions introduced into the TOF mass analyzer are accelerated along the axis perpendicular to their initial direction of motion.



Tandem time-of-flight (TOF/TOF)

A tandem mass spectrometry method where two time-of-flight mass spectrometers are used consecutively.



Properties of mass spectrometer

- Mass Resolution
- Mass accuracy
- Sensitivity
- Duty cycle
- Throughput

Mass Resolution

Resolution

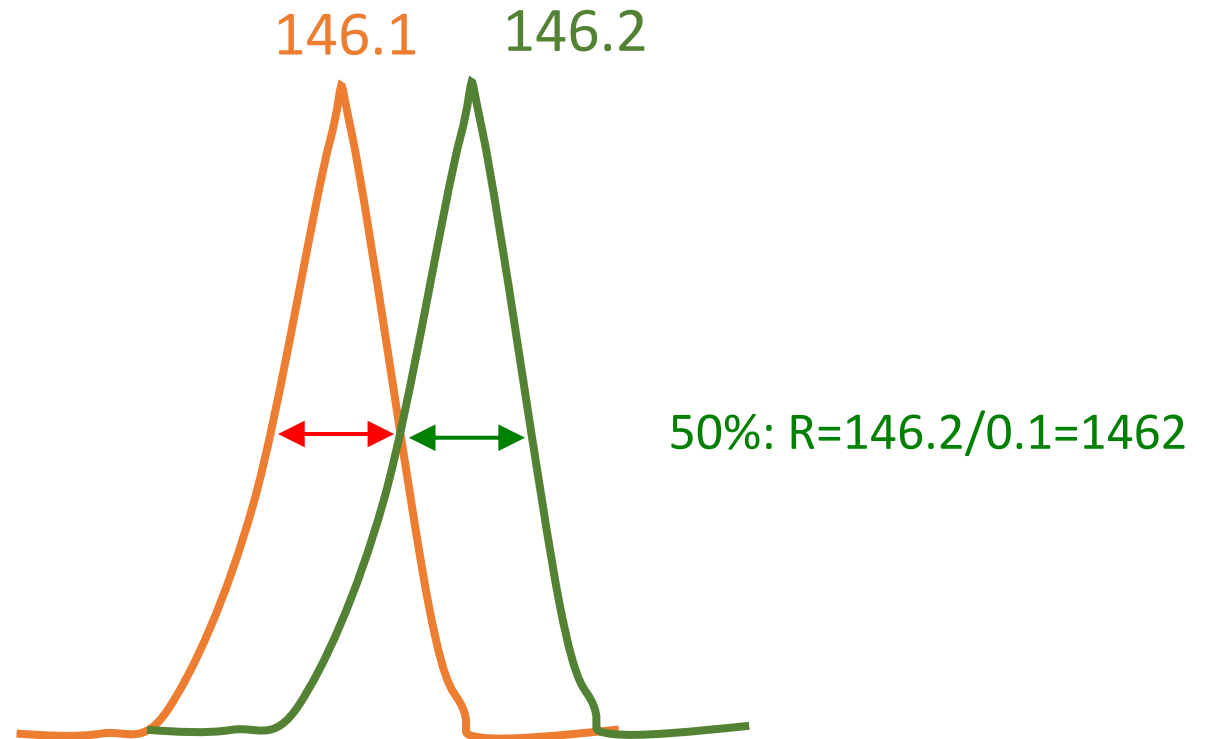
$$R = M/\Delta M$$

R: resolving power

M: mass of the (second) peak

ΔM : Peak width

Resolution may change with M

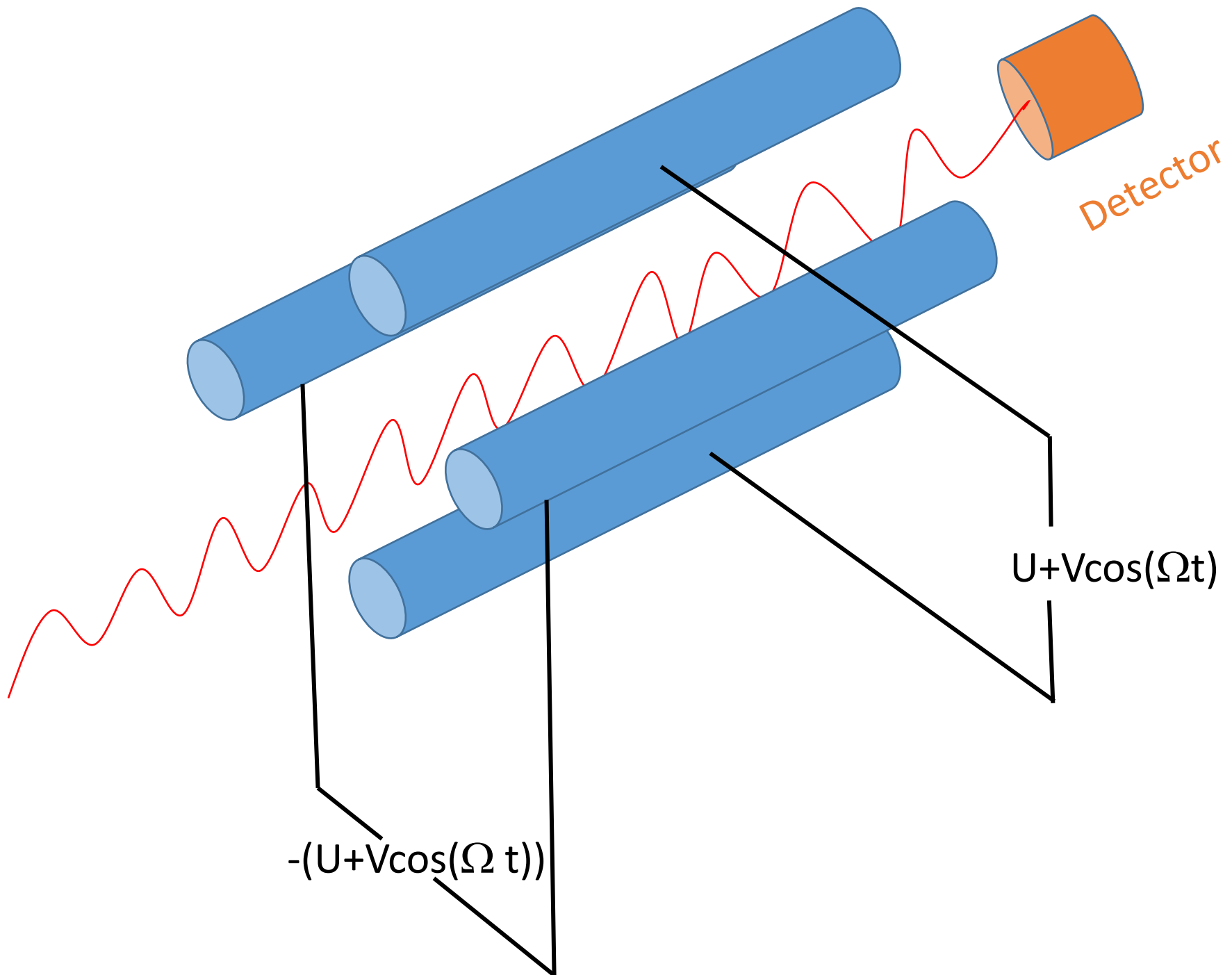


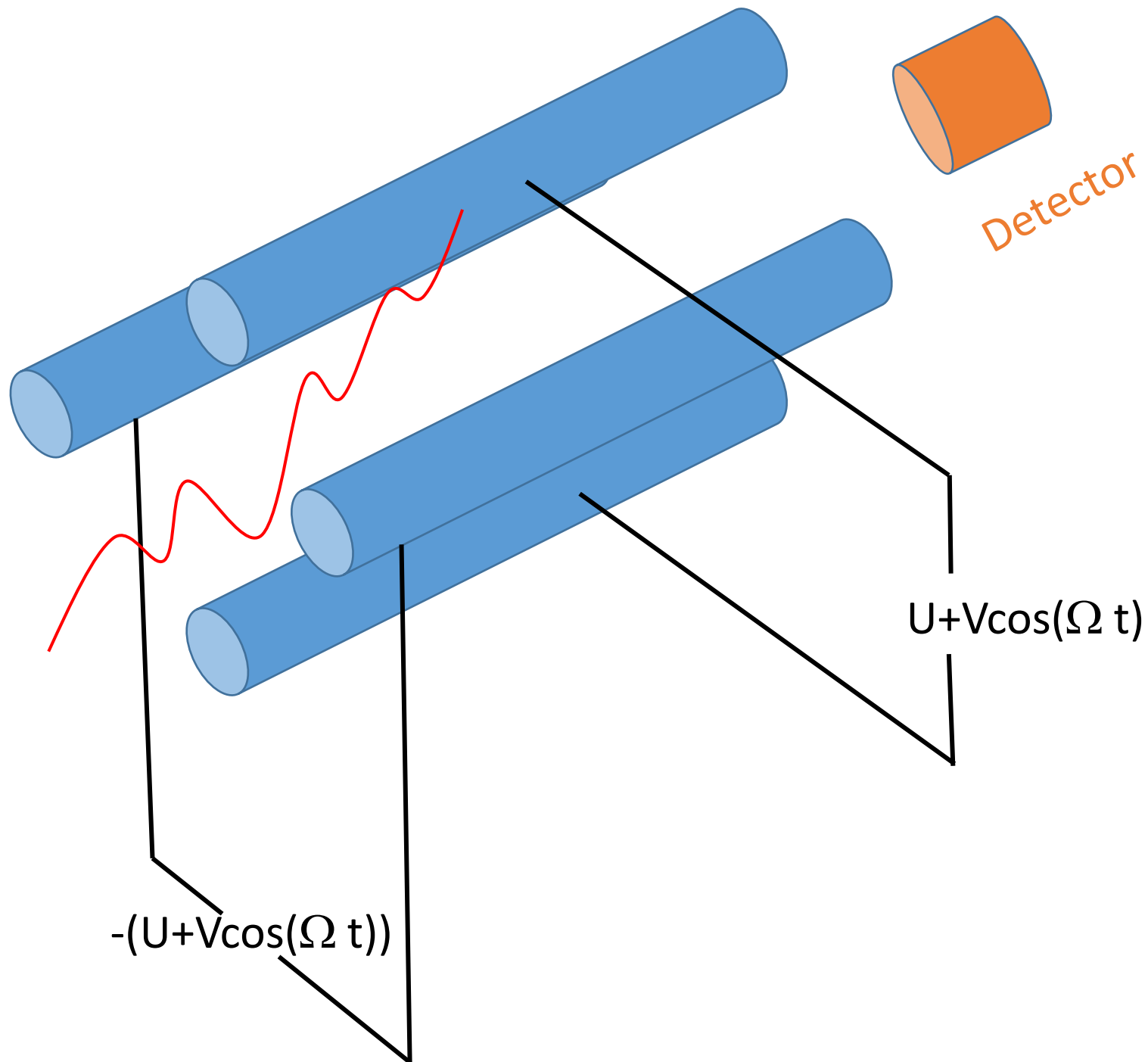
Quadrupole mass spectrometer

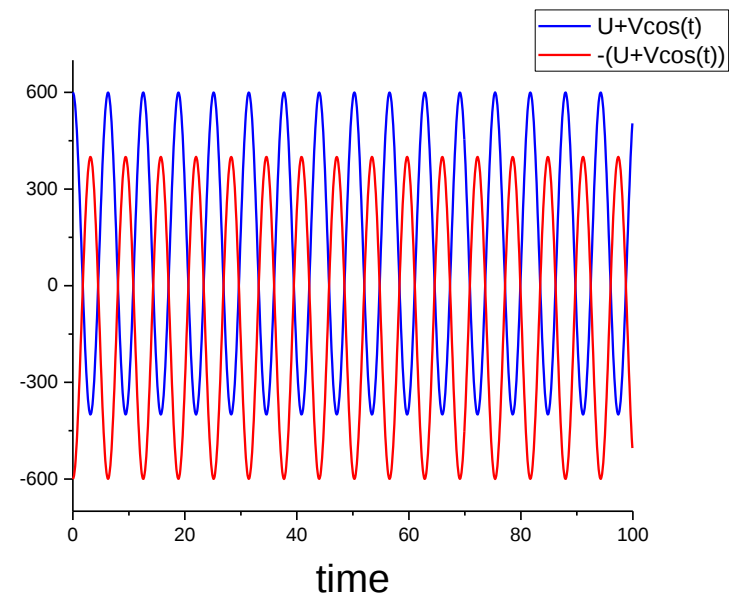
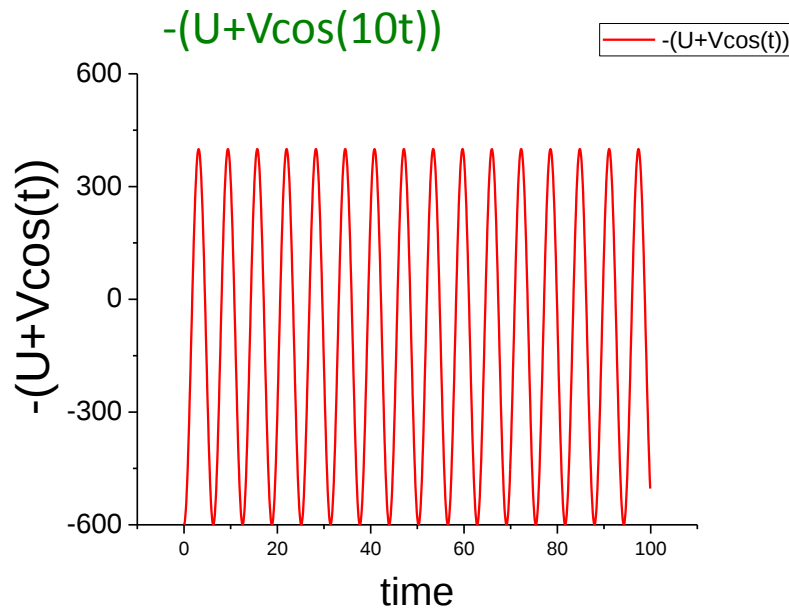
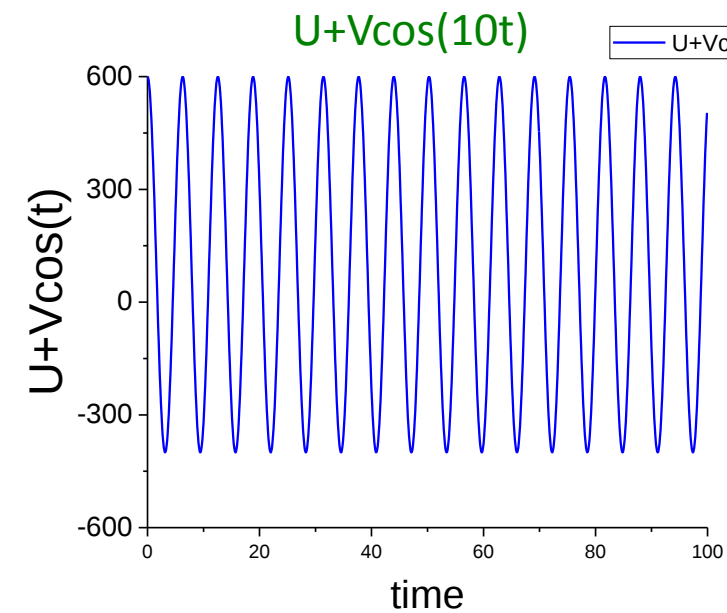
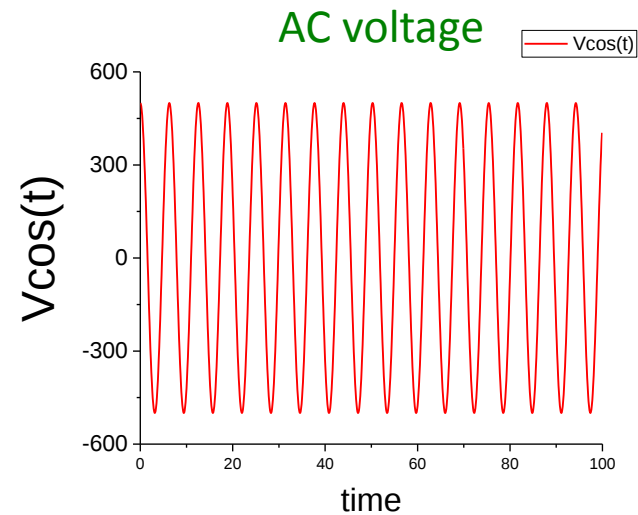
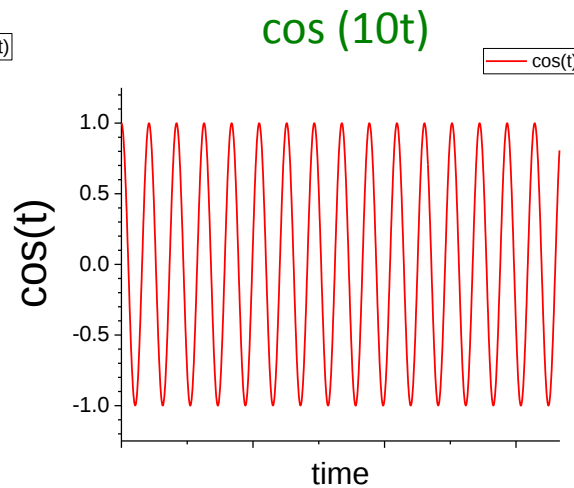
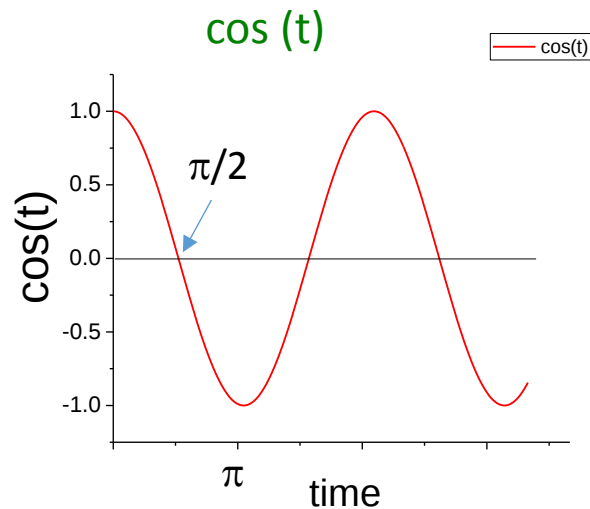
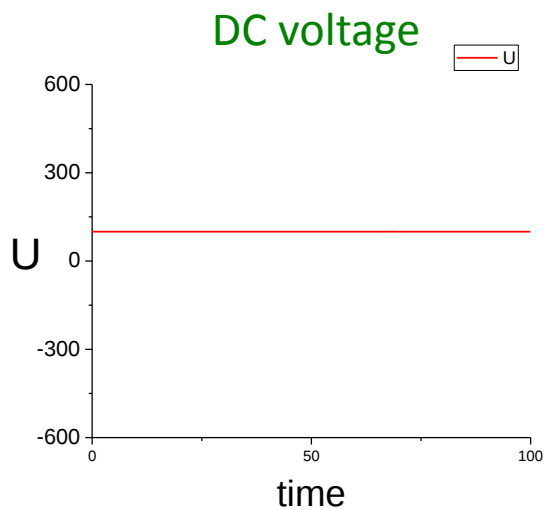


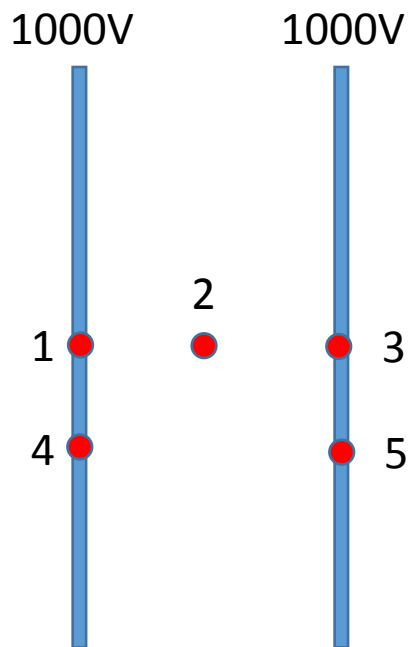
Wolfgang Paul

1989 Nobel prize in Physics



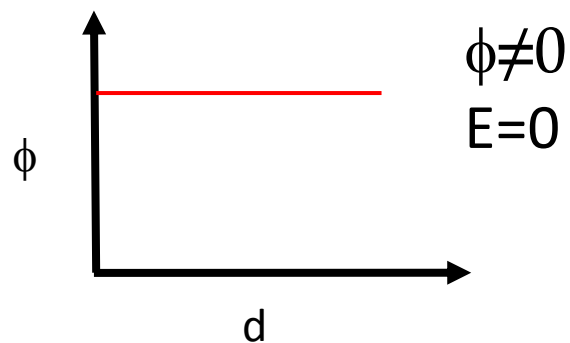




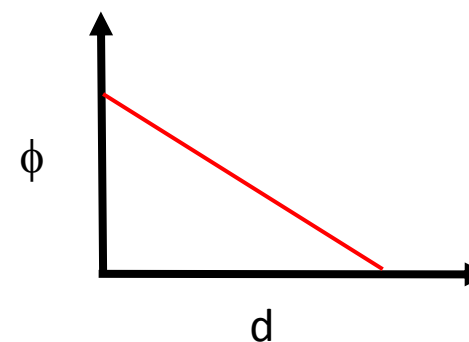
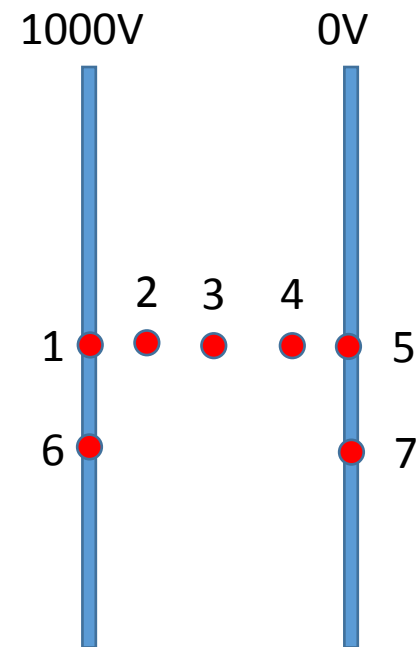


Potential ϕ
&
Electric field E

$$E = -\frac{\partial \phi}{\partial d}$$

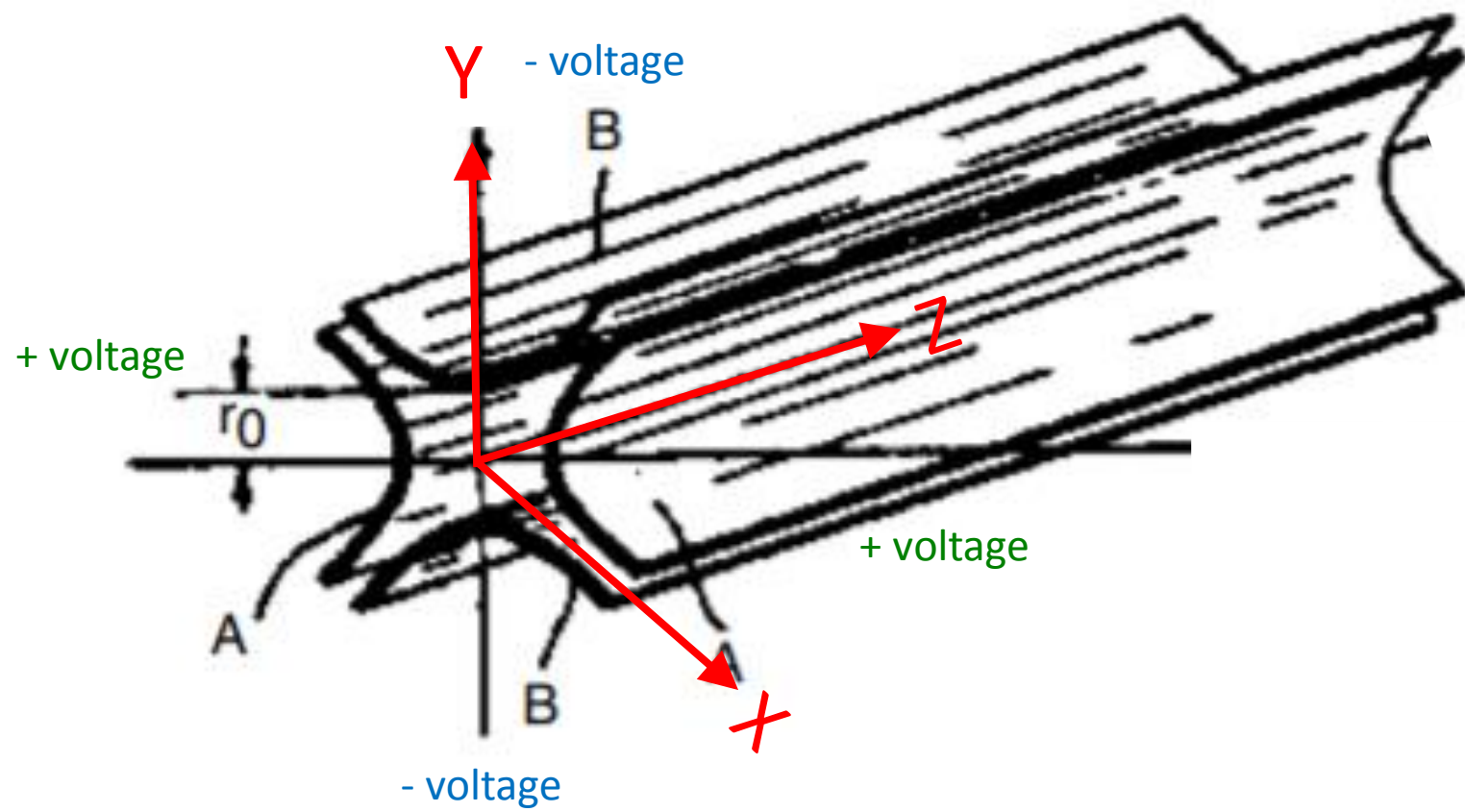


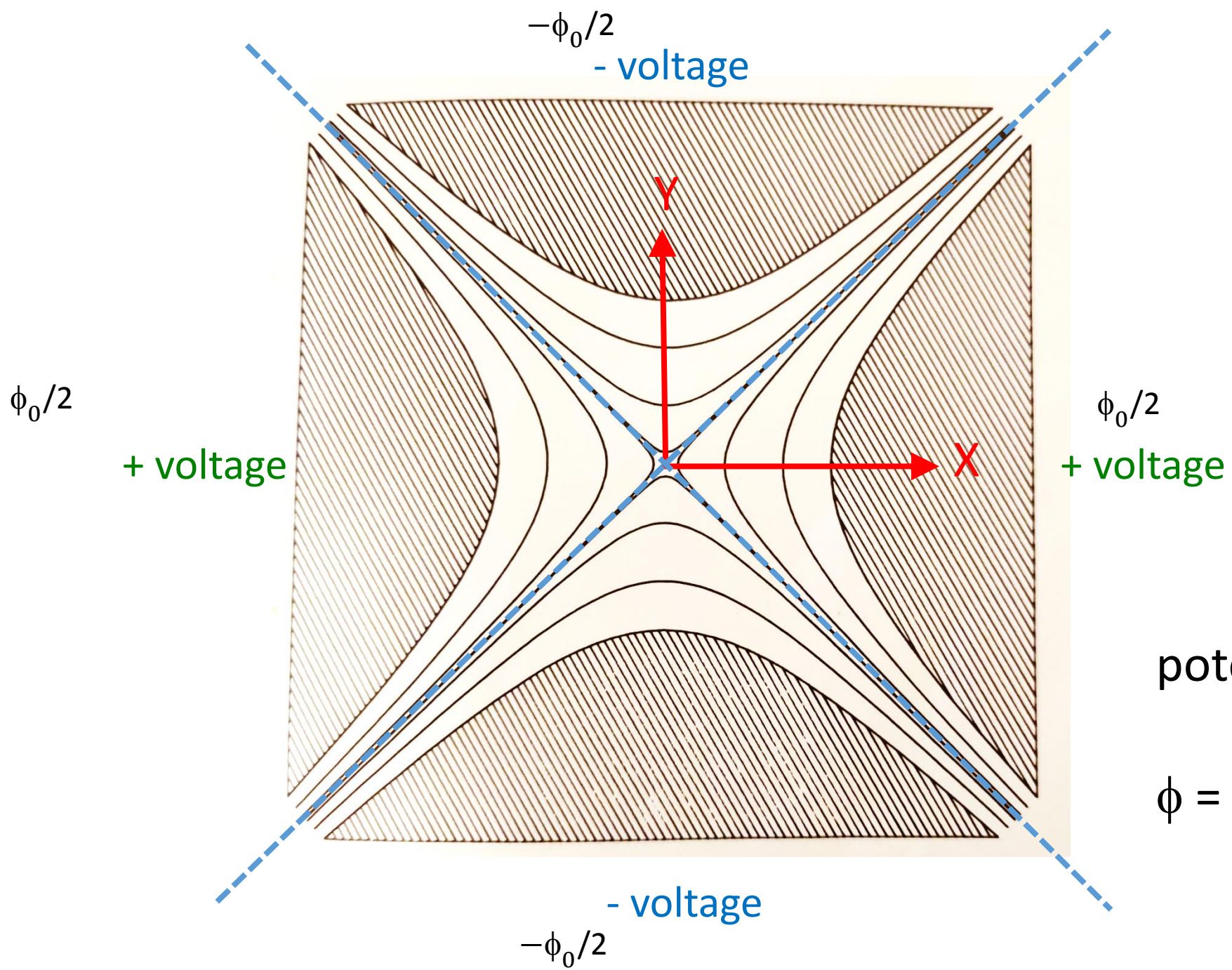
$$\phi_1 = \phi_2 = \phi_3 = \phi_4 = \phi_5 = 1000 \text{ V}$$



$$\phi_1 = \phi_6 = 1000 \text{ V} \qquad \phi_5 = \phi_7 = 0 \text{ V}$$

$$\phi_2 = 750 \text{ V} \quad \phi_3 = 500 \text{ V} \quad \phi_4 = 250 \text{ V}$$





$$F_x = -eE_x = -e \frac{\partial \phi}{\partial x} = -e \frac{\phi_0}{2r_0^2} (2x) \quad (1)$$

potential

$$F_y = -eE_y = -e \frac{\partial \phi}{\partial y} = -e \frac{\phi_0}{2r_0^2} (-2y) \quad (2)$$

$$\phi = \frac{\phi_0}{2r_0^2} (x^2 - y^2)$$

$$F_x = ma_x = m \frac{\partial^2 x}{\partial t^2} \quad (3)$$

$$F_y = ma_y = m \frac{\partial^2 y}{\partial t^2} \quad (4)$$

From (1) and (3):

$$m \frac{\partial^2 x}{\partial t^2} = -e \frac{\phi_0}{2r_0^2} (2x)$$

From (2) and (4):

$$m \frac{\partial^2 y}{\partial t^2} = -e \frac{\phi_0}{2r_0^2} (-2y)$$

Motions in x and y directions
are independent

What is ϕ_0 ?

(1) $\phi_0 = 2(U + V\cos(\Omega t))$ Conventional method

(2) $\phi_0 = \text{others}$ e.g., digital ion trap

For conventional method:

$$m \frac{\partial^2 x}{\partial t^2} = -e \frac{\phi_0}{r_0^2} (2x) = -e \frac{2(U + V\cos(\Omega t))}{2r_0^2} (2x)$$

$$m \frac{\partial^2 y}{\partial t^2} = -e \frac{\phi_0}{r_0^2} (-2y) = -e \frac{2(U + V\cos(\Omega t))}{2r_0^2} (-2y)$$

For conventional method:

$$m \frac{\partial^2 x}{\partial t^2} = -e \frac{\phi_0}{r_0^2} (2x) = -e \frac{2(U + V \cos(\Omega t))}{2r_0^2} (2x)$$

$$\frac{\partial^2 x}{\partial t^2} = -\frac{2e}{mr_0^2} (U + V \cos(\Omega t)) x$$

$$m \frac{\partial^2 y}{\partial t^2} = -e \frac{\phi_0}{r_0^2} (-2y) = -e \frac{2(U + V \cos(\Omega t))}{2r_0^2} (-2y)$$

$$\frac{\partial^2 y}{\partial t^2} = \frac{2e}{mr_0^2} (U + V \cos(\Omega t)) y$$

$$\frac{\partial^2 x}{\partial t^2} = -\frac{2e}{mr_0^2} (U + V \cos(\Omega t)) x$$

$$\frac{\partial^2 y}{\partial t^2} = \frac{2e}{mr_0^2} (U + V \cos(\Omega t)) y$$

$$\frac{\partial^2 x}{\partial t^2} + \left(\frac{2e}{mr_0^2} U + \frac{2e}{mr_0^2} V \cos(\Omega t) \right) x = 0$$

$$\frac{\partial^2 y}{\partial t^2} - \left(\frac{2e}{mr_0^2} U + \frac{2e}{mr_0^2} V \cos(\Omega t) \right) y = 0$$

$$\xi = \Omega t / 2$$

$$u = x \text{ or } y$$

The Mathieu Equation

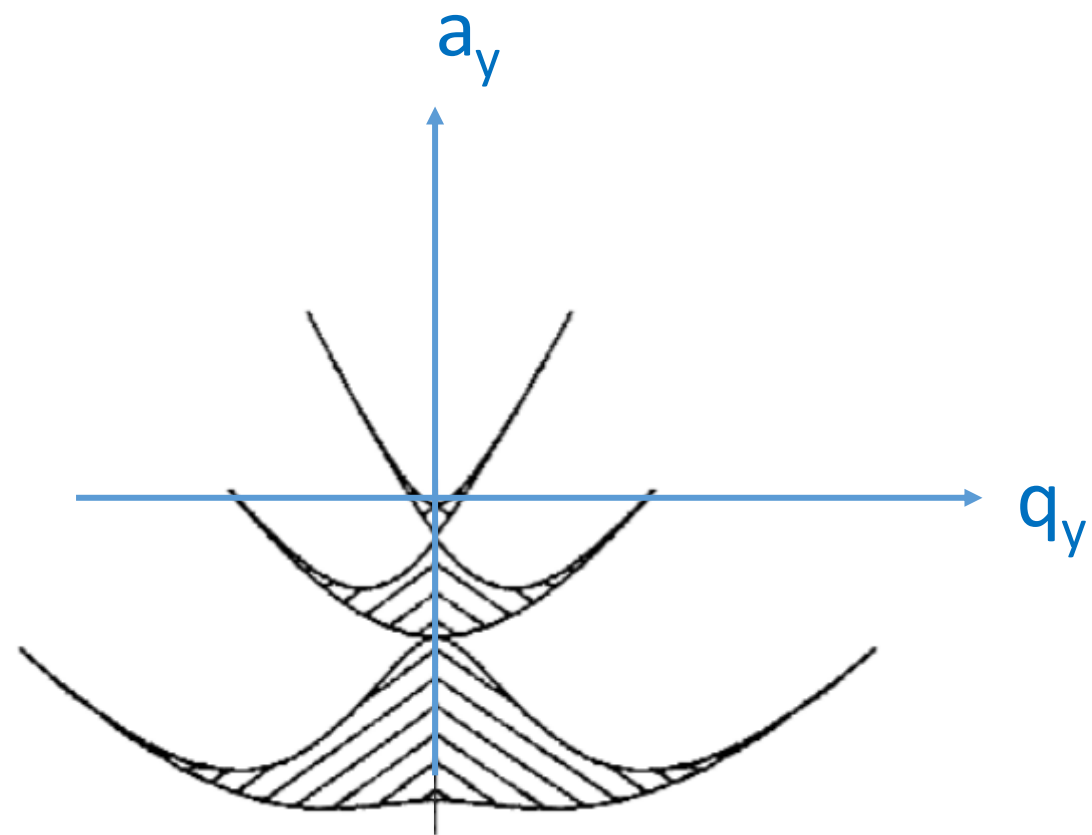
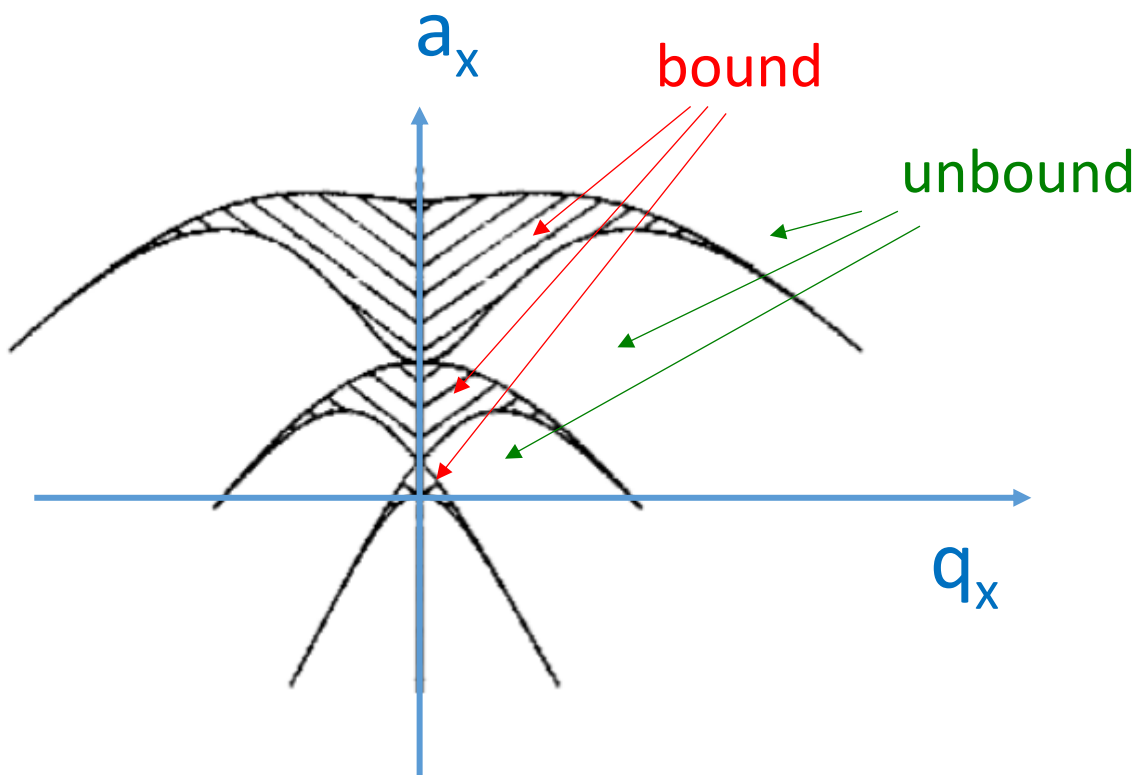
$$\frac{d^2 u}{d\xi^2} + (a_u + 2q_u \cos 2\xi) u = 0$$

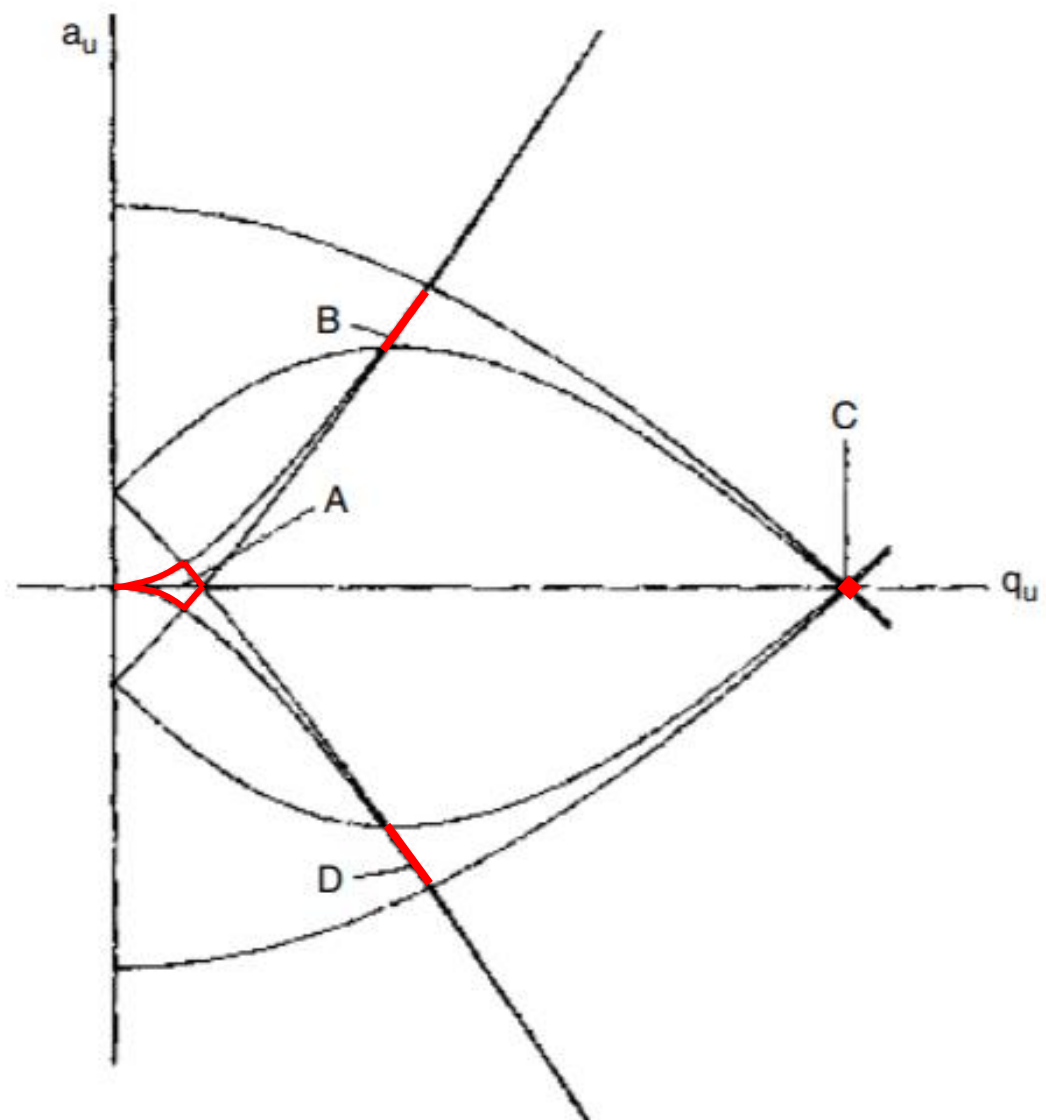
$$a_x = \frac{8eU}{mr_0^2 \Omega^2}$$

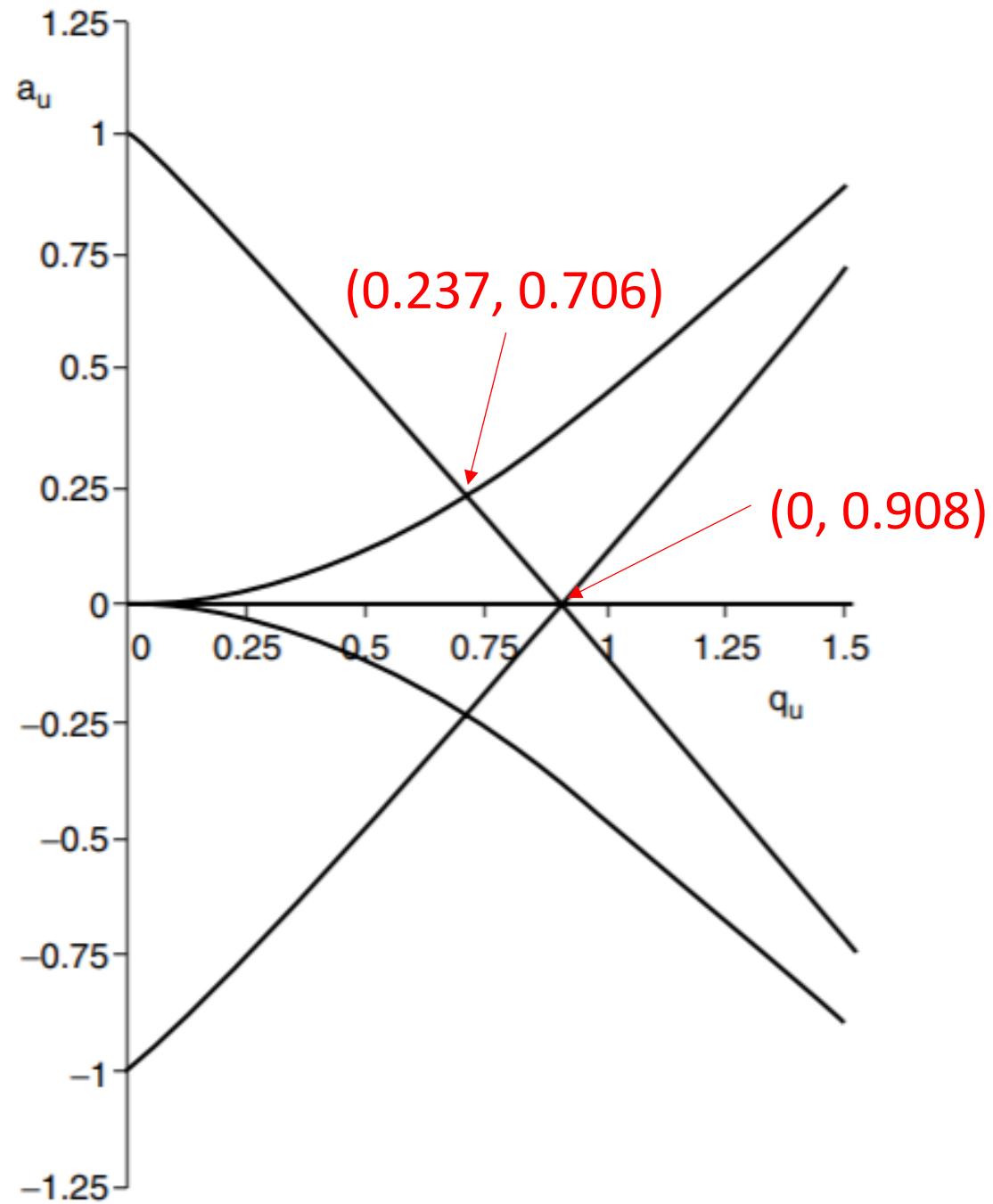
$$q_x = \frac{4eV}{mr_0^2 \Omega^2}$$

$$a_y = \frac{-8eU}{mr_0^2 \Omega^2}$$

$$q_y = \frac{-4eV}{mr_0^2 \Omega^2}$$





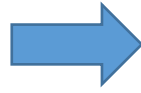


$$a_u = \frac{8eU}{mr_0^2\Omega^2}$$

$$q_u = \frac{4eV}{mr_0^2\Omega^2}$$

$$a_x = \frac{8eU}{mr_0^2\Omega^2}$$

$$q_x = \frac{4eV}{mr_0^2\Omega^2}$$



$$U = \frac{mr_0^2\Omega^2}{8e} a_x$$

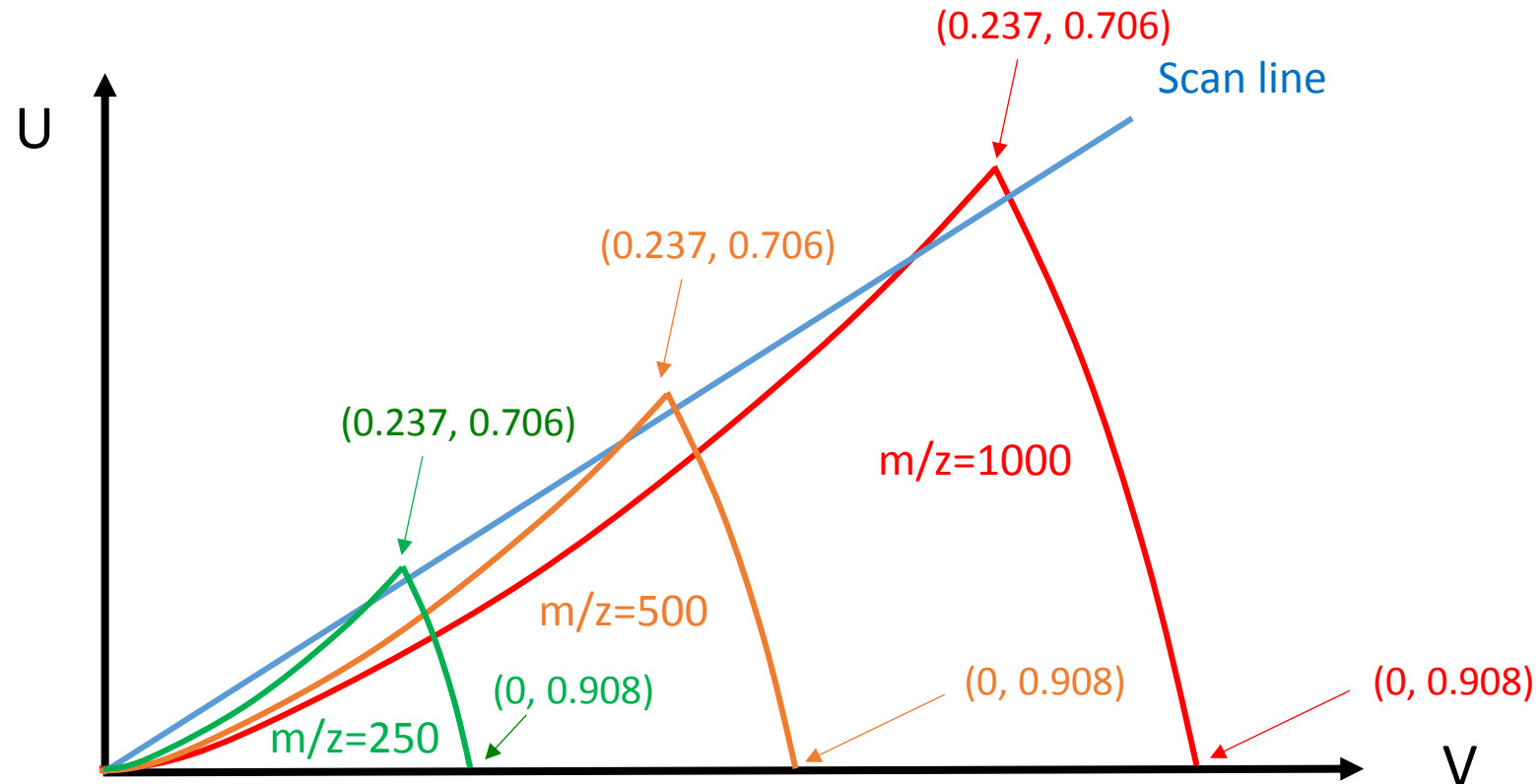
$$V = \frac{mr_0^2\Omega^2}{4e} q_x$$

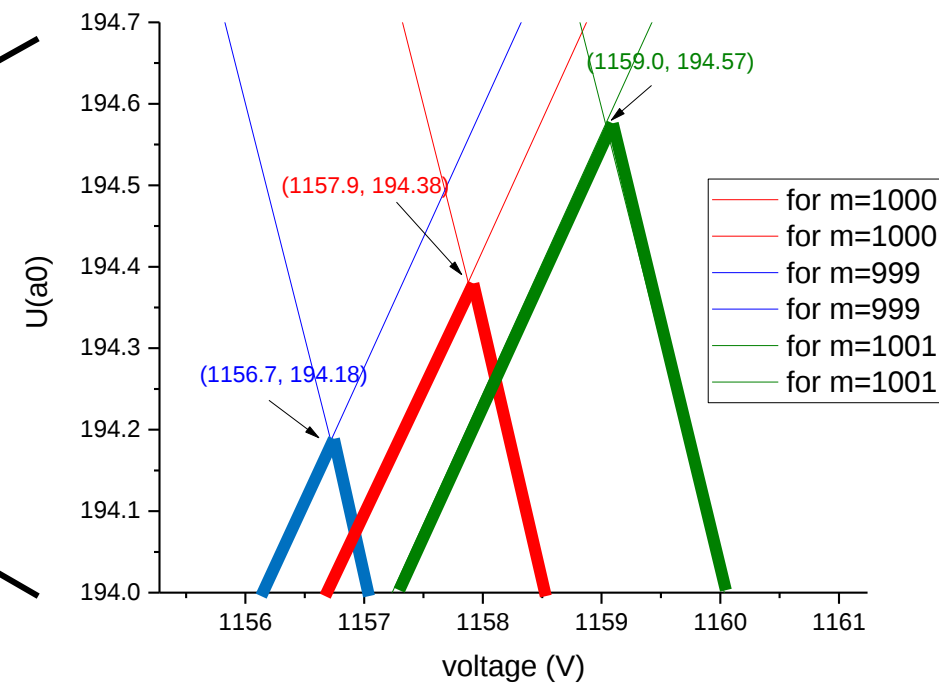
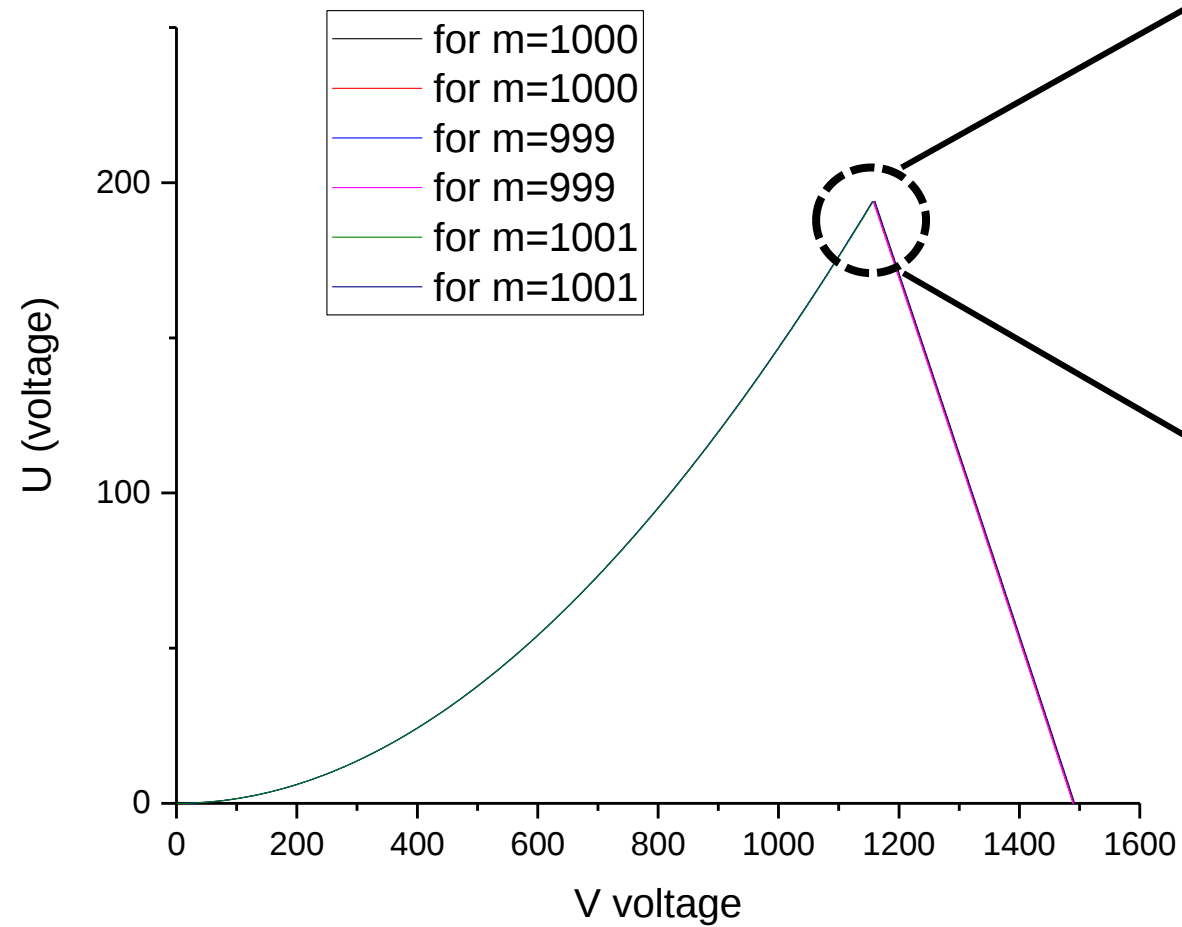
For a given r_0 and Ω : e.g., $r_0 = 4 \text{ mm}$ and $\Omega = 1\text{MHz}$

$$V = \frac{1000\text{g/mol} \times 10^{-3}\text{kg/g} \times (0.004\text{m})^2 (2\pi \times 1 \times 10^6 \text{rad/s})^2}{6.02 \times 10^{23} \text{ 1/mol} \times 4 \times 1.6 \times 10^{-19} \text{C}} \times 0.706 = 1157.4 \text{ V (V}_{0\text{-p}})$$

Mass upper limit: available voltage

$$U = \frac{1000\text{g/mol} \times 10^{-3}\text{kg/g} \times (0.004\text{m})^2 (2\pi \times 1 \times 10^6 \text{rad/s})^2}{6.02 \times 10^{23} \text{ 1/mol} \times 8 \times 1.6 \times 10^{-19} \text{C}} \times 0.237 = 194.2 \text{ V (DC)}$$





At apex:

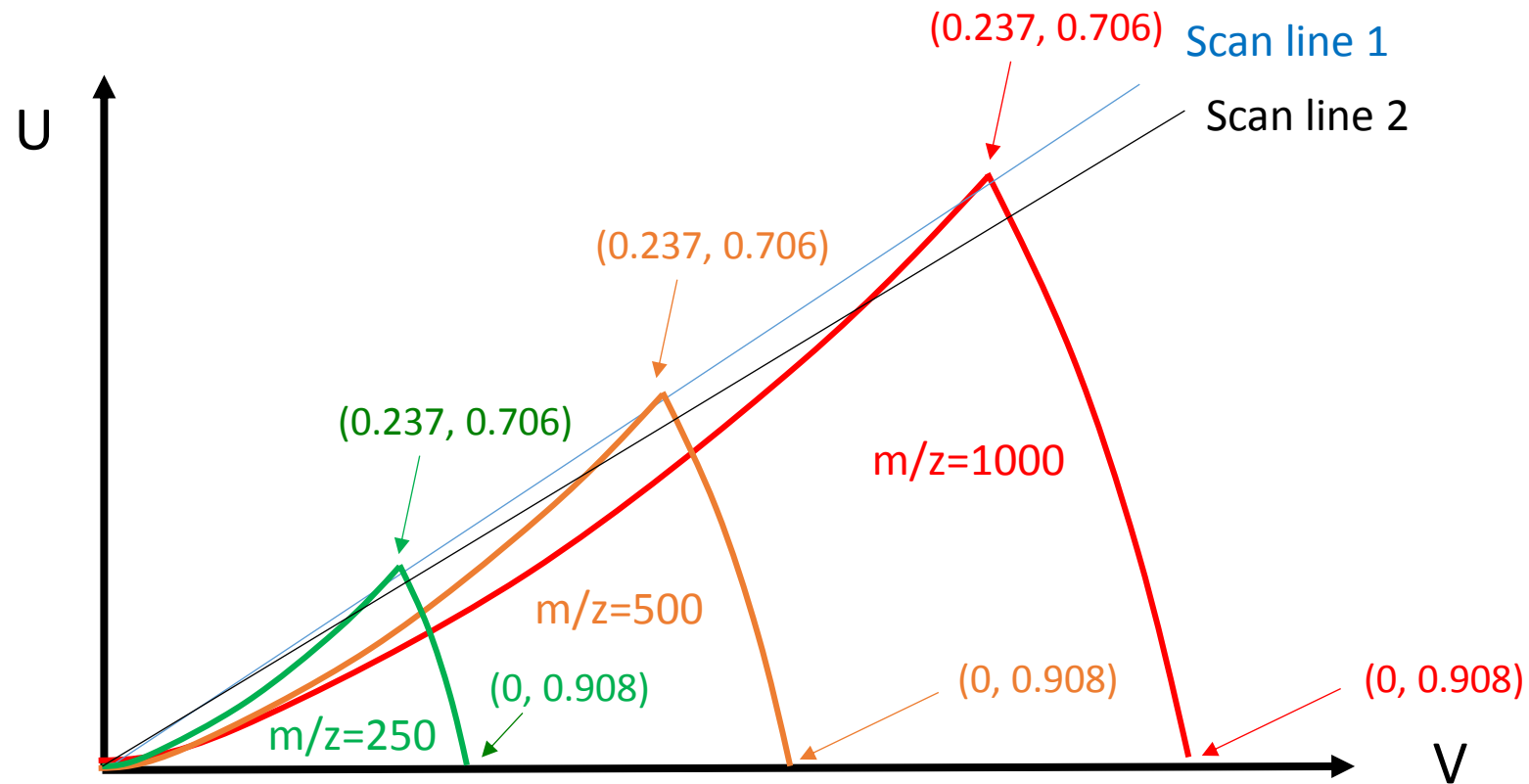
$$\Delta U = \frac{m_1 r_0^2 \Omega^2}{8e} a_x - \frac{m_2 r_0^2 \Omega^2}{8e} a_x = \Delta m \frac{r_0^2 \Omega^2}{8e} a_x = \frac{1 \times 10^{-3}}{6.02 \times 10^{23}} \times \frac{0.004^2 (2\pi \times 10^6)^2}{8 \times 1.6 \times 10^{-19}} 0.237 = 0.19 \text{ V}$$

$$\Delta V = \frac{m_1 r_0^2 \Omega^2}{4e} q_x - \frac{m_2 r_0^2 \Omega^2}{4e} q_x = \Delta m \frac{r_0^2 \Omega^2}{4e} q_x = \frac{1 \times 10^{-3}}{6.02 \times 10^{23}} \times \frac{0.004^2 (2\pi \times 10^6)^2}{4 \times 1.6 \times 10^{-19}} 0.706 = 1.13 \text{ V}$$

$\Delta U, \Delta V$ are m independent

For $m/\Delta m > 1000$

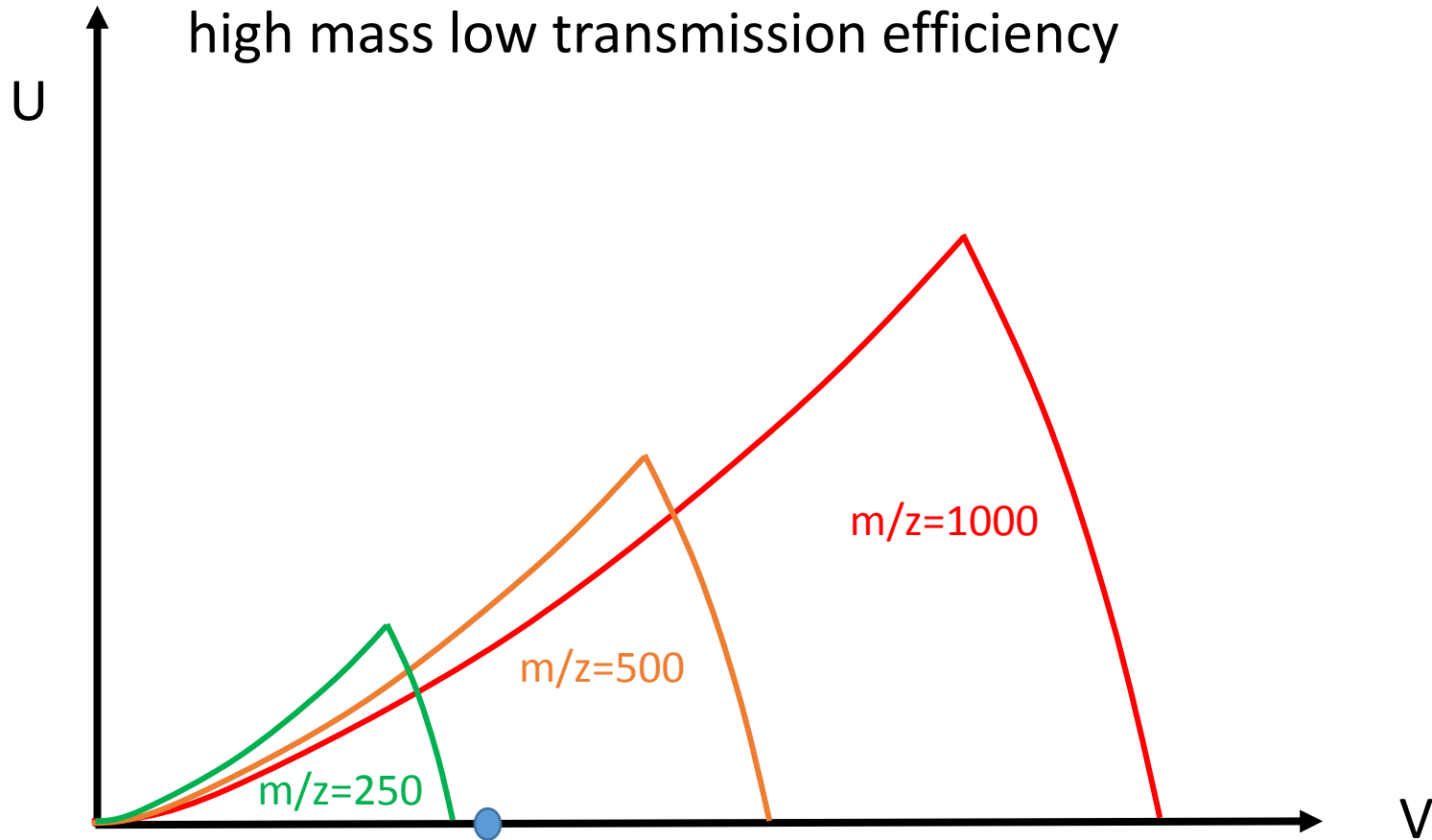
- $\Delta\Omega/\Omega < 1/2000$
- $\Delta r_0/r_0 < 1/2000$



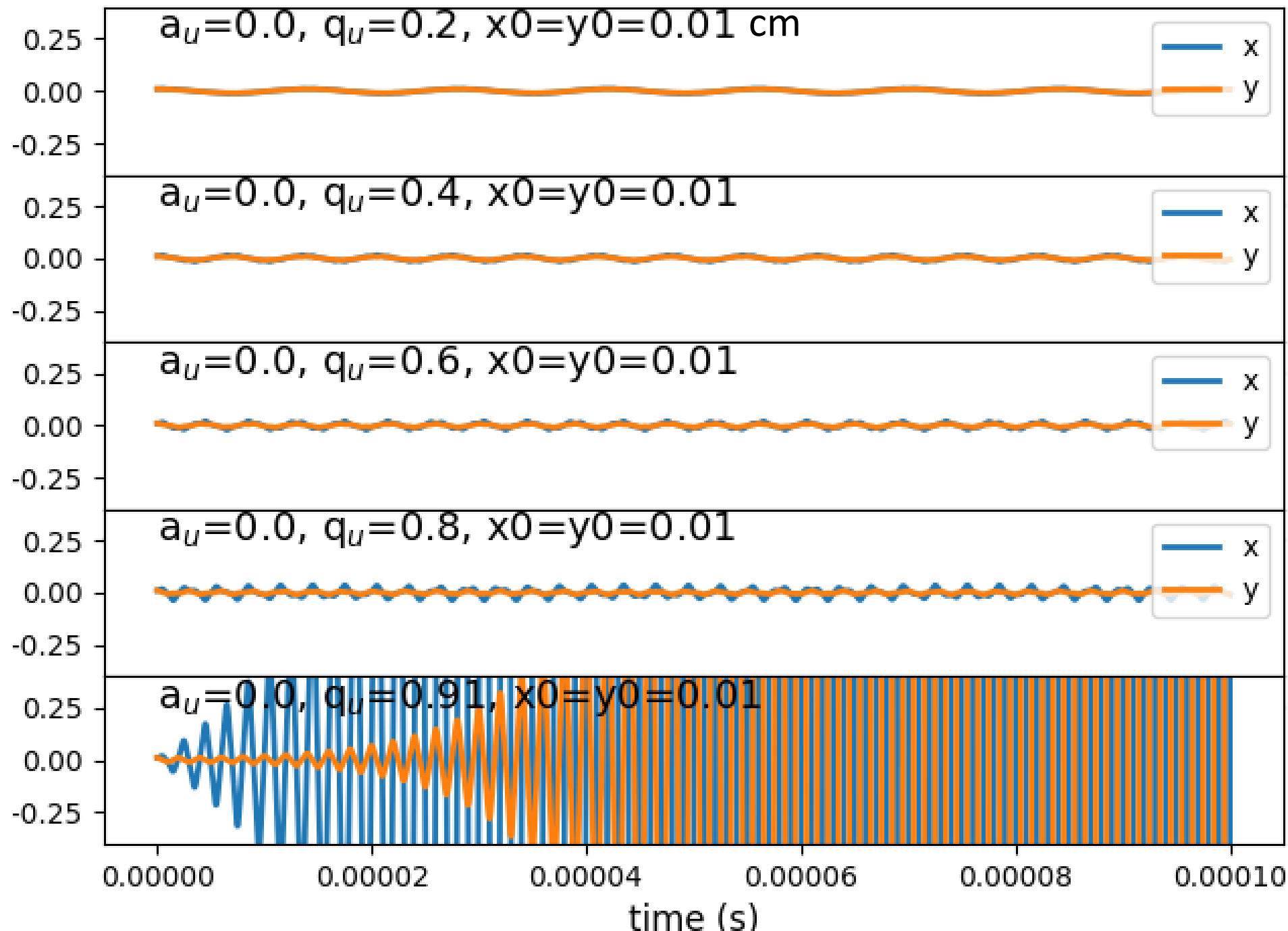
Scan line 1: $U/V = 0.5 \times (0.237/0.706) = 0.1678$, best resolution, small signal

Scan line 2: $U/V < 0.5 \times (0.237/0.706) = 0.1678$, low resolution, large signal

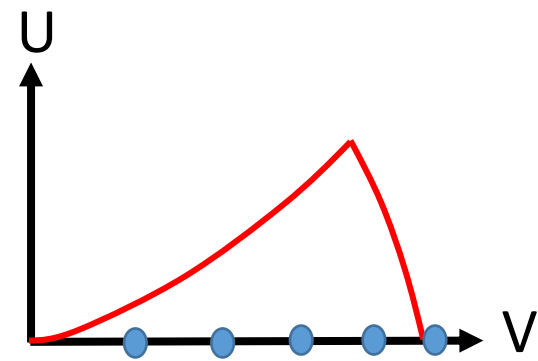
Ion guide: turn off DC voltage ($U=0$), $a_x=0$
low mass cutoff
high mass low transmission efficiency

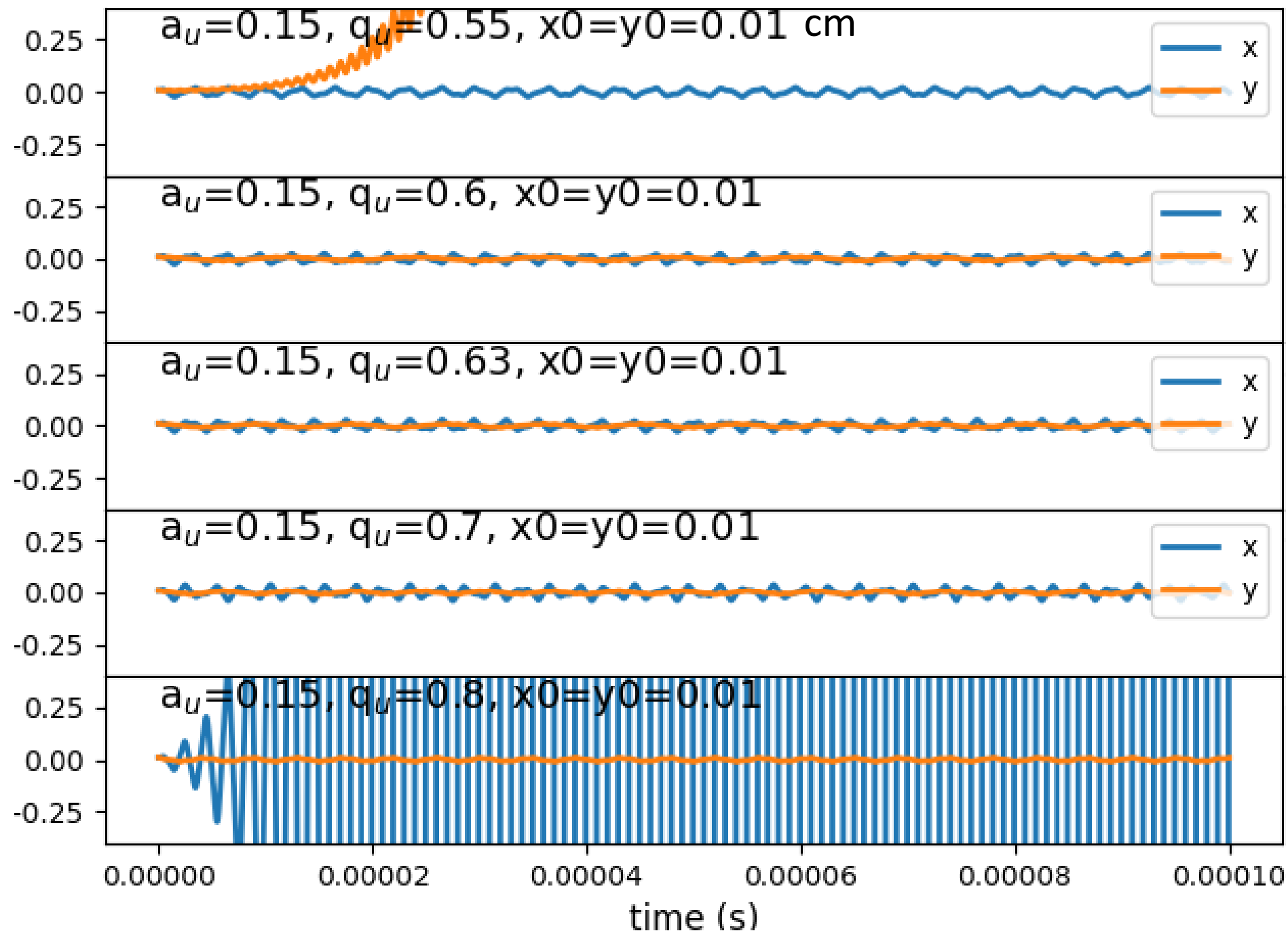


Trajectory

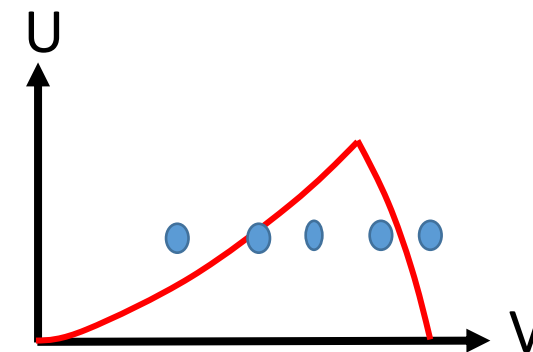


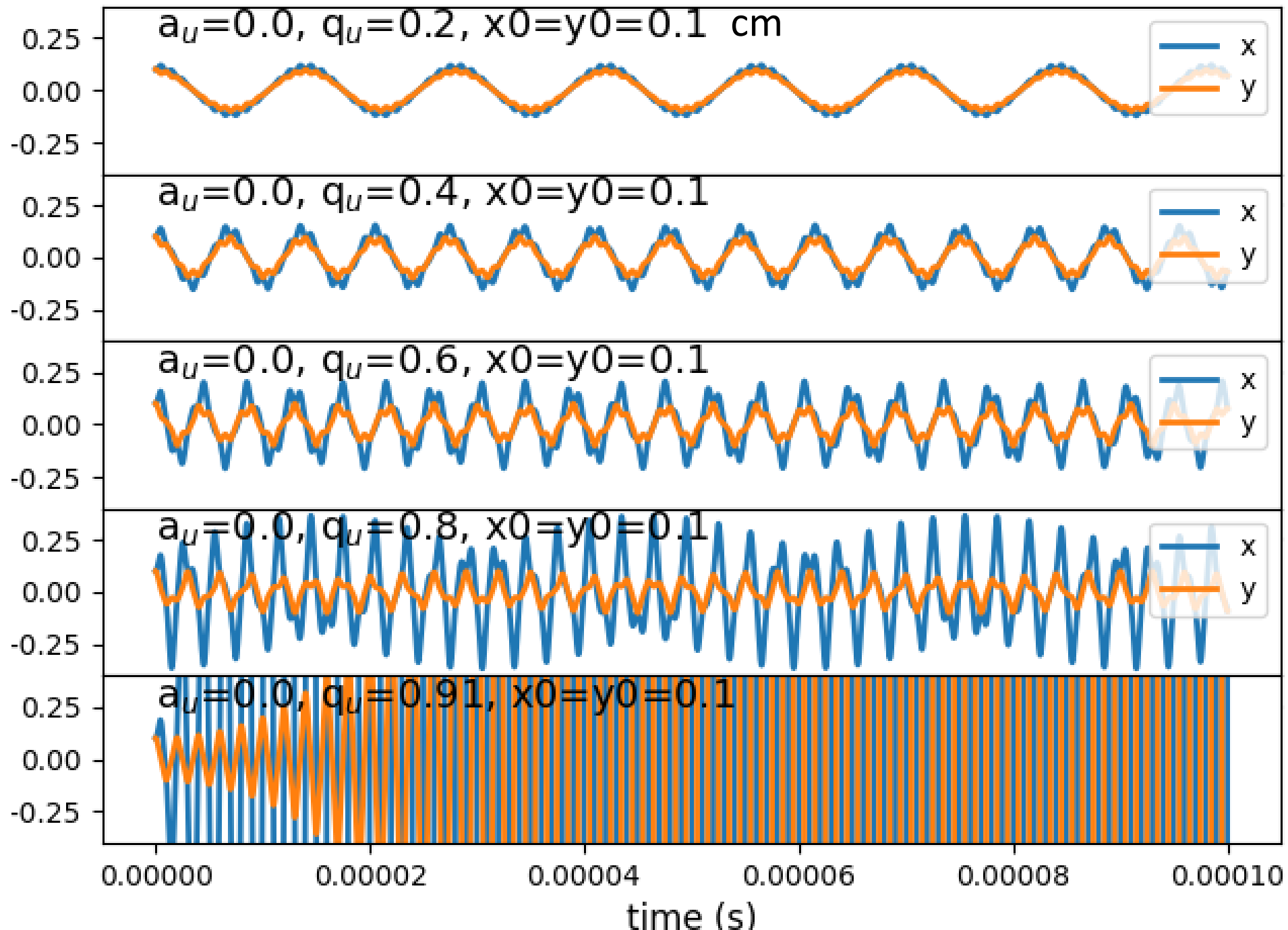
Initial velocity
 $V_x = V_y = 0$
Phase = 0



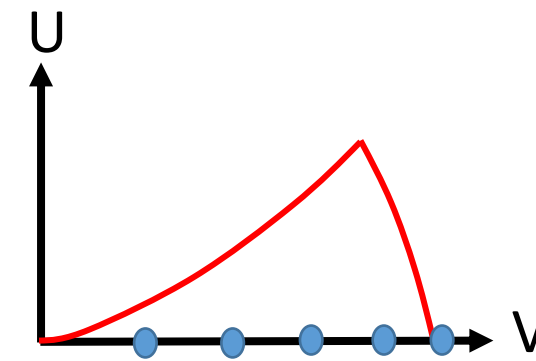


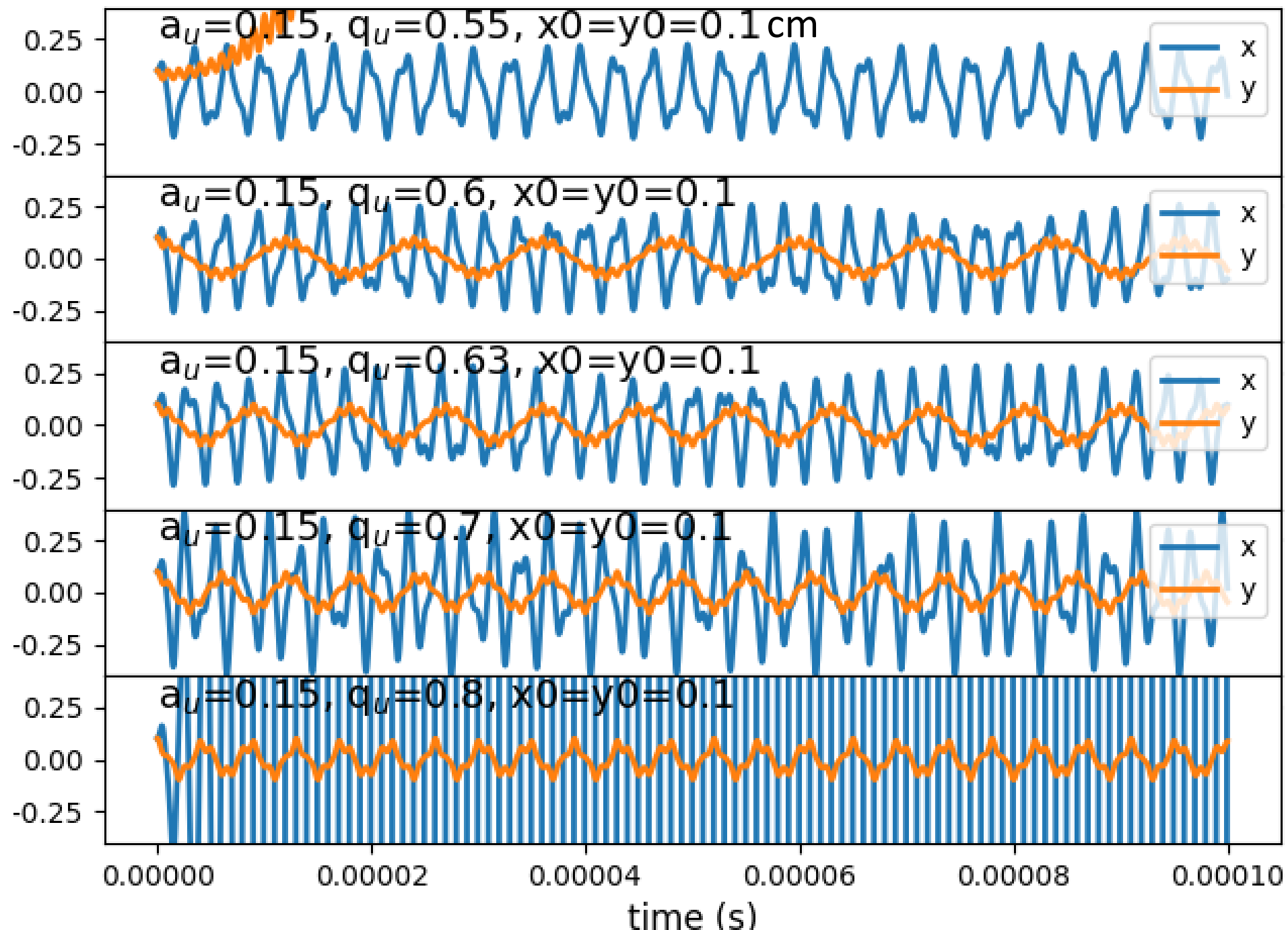
Initial velocity
 $V_x = V_y = 0$
Phase = 0





Initial velocity
 $V_x = V_y = 0$
Phase = 0





Initial velocity
 $V_x = V_y = 0$
Phase = 0

